# Adiabatic and Nonadiabatic Reaction Pathways of the $\mathbf{O}\left({ }^{3} \mathbf{P}\right)$ with Propyne 

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#### Abstract

For the reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with propyne, the product channels and mechanisms are investigated both theoretically and experimentally. Theoretically, the $\operatorname{CCSD}(\mathrm{T}) / / \mathrm{B} 3 \mathrm{LYP} / 6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level of calculations are performed for both the triplet and singlet potential energy surfaces and the minimum energy crossing point between the two surfaces are located with the Newton-Lagrange method. The theoretical calculations show that the reaction occurs dominantly via the O -addition rather than the H -abstraction mechanism. The reaction starts with the O-addition to either of the triple bond carbon atoms forming triplet ketocarbene ${ }^{3} \mathrm{CH}_{3} \mathrm{CCHO}$ or ${ }^{3} \mathrm{CH}_{3} \mathrm{COCH}$ which can undergo decomposition, H -atom migration or intersystem crossing from which a variety of channels are open, including the adiabatic channels of $\mathrm{CH}_{3} \mathrm{CCO}+\mathrm{H}\left(\mathrm{CH}_{2} \mathrm{CCHO}+\mathrm{H}\right), \mathrm{CH}_{3}+\mathrm{HCCO}, \mathrm{CH}_{2} \mathrm{CH}+$ $\mathrm{HCO}, \mathrm{CH}_{2} \mathrm{CO}+\mathrm{CH}_{2}, \mathrm{CH}_{3} \mathrm{CH}+\mathrm{CO}$, and the nonadiabatic channels of $\mathrm{C}_{2} \mathrm{H}_{4}+\mathrm{CO}, \mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{H}_{2}+\mathrm{CO}, \mathrm{H}_{2}$ $+\mathrm{H}_{2} \mathrm{CCCO}$. Experimentally, the CO channel is investigated with TR-FTIR emission spectroscopy. A complete detection of the CO product at each vibrationally excited level up to $v=5$ is fulfilled, from which the vibrational energy disposal of CO is determined and found to consist with the statistical partition of the singlet $\mathrm{C}_{2} \mathrm{H}_{4}+\mathrm{CO}$ channel, but not with the triplet $\mathrm{CH}_{3} \mathrm{CH}+\mathrm{CO}$ channel. In combination with the present calculation results, it is concluded that CO arises mainly from the singlet methylketene $\left({ }^{1} \mathrm{CH}_{3} \mathrm{CHCO}\right)$ dissociation following the intersystem crossing of the triplet ketocarbene adduct ( $\left.{ }^{3} \mathrm{CH}_{3} \mathrm{CCHO}\right)$. Fast intersystem crossing via the minimum energy crossing point of the triplet and singlet surfaces is shown to play significant roles resulting into nonadiabatic pathways for this reaction. Moreover, other interesting questions are explored as to the site selectivity of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ atom being added to which carbon atom of the triple bond and different types of internal H -atom migrations including $1,2-\mathrm{H}$ shift, $3,2-\mathrm{H}$ shift, and $3,1-\mathrm{H}$ shift involved in the reaction.


## 1. Introduction

The actual oxidation process of hydrocarbon fuels often involves the reactions of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with unsaturated hydrocarbons including alkenes and alkynes. The reactions of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with alkynes are therefore of importance in combustion. Also the addition reactions of the electrophilic oxygen atoms to the triple bond of alkynes serve as fascinating examples for chemical dynamics studies. The reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with alkynes has been attracting much attention of both experimental and theoretical studies. ${ }^{1-22}$

The reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with the simplest alkyne, $\mathrm{C}_{2} \mathrm{H}_{2}$, has been studied extensively. Experimental ${ }^{1-12}$ and theoretical ${ }^{13-15}$ studies agree that the primary products of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{C}_{2} \mathrm{H}_{2}$ reaction are produced mainly through two channels 1 a and 1 b . However, the potential H -atom abstraction channel, 1c, is negligible because of its large endothermicity of $31 \mathrm{kcal} \mathrm{mol}^{-1} .5,12,15$
$\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{HCCO}+\mathrm{H} \quad \Delta H=-19 \mathrm{kcal} \mathrm{mol}^{-1}$
$\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{CH}_{2}+\mathrm{CO} \quad \Delta H=-47 \mathrm{kcal} \mathrm{mol}^{-1}$
$\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{OH}+\mathrm{CCH} \quad \Delta H=+31 \mathrm{kcal} \mathrm{mol}^{-1}$
It has been well established that the major channel of the $\mathrm{C}_{2} \mathrm{H}_{2}$ $+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ reaction is eq 1 a , the $\mathrm{HCCO}+\mathrm{H}$ channel. ${ }^{1,3,5-7,9}$ Recent experiments ${ }^{5-7,9}$ determined the yield of the $\mathrm{HCCO}+\mathrm{H}$ channel to be about $80 \%$, and the $\mathrm{CH}_{2}+\mathrm{CO}$ channel to be $15-20 \%$.

[^0]The calculation of Nguyen et al. ${ }^{15}$ suggested that the HCCO + H yield is $93 \%$ at 300 K and drops to $90 \%$ at 1000 K , whereas the $\mathrm{CH}_{2}+\mathrm{CO}$ yield rises from $7 \%$ at 300 K to $10 \%$ at 1000 K . Their computation also shows that the $\mathrm{HCCO}+\mathrm{H}$ yield dominates over the $\mathrm{CH}_{2}+\mathrm{CO}$ yield.

The mechanism of the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{C}_{2} \mathrm{H}_{2}$ reaction can be characterized as the electrophilic ground-state oxygen atoms being added to $\mathrm{C}-\mathrm{C}$ triple bond, ${ }^{12}$ forming a diradical intermediate HCCHO which is a ketocarbene initially. This diradical intermediate or its isomer, ketene $\mathrm{CH}_{2} \mathrm{CO}$ formed via $1,2-\mathrm{H}$ shift, can subsequently undergo unimolecular decomposition into $\mathrm{HCCO}+\mathrm{H}$ (1a) or $\mathrm{CH}_{2}+\mathrm{CO}$ (1b), respectively. Recently, the potential energy surfaces for the $\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ reaction were theoretically investigated in detail. ${ }^{15}$ Except for the two major channels presented above, the H -abstraction channel was calculated to have a high endothermicity ( $+31.9 \mathrm{kcal} \mathrm{mol}^{-1}$ ) and thus cannot compete with the addition/elimination channels under any combustion conditions. In addition, an efficient reaction pathway on the electronically excited ${ }^{3} \mathrm{~A}^{\prime}$ surface resulting into $\mathrm{H}\left({ }^{2} \mathrm{~S}\right)+\mathrm{HCCO}\left(\mathrm{A}^{2} \mathrm{~A}^{\prime}\right)$ is predicted to play an important role at higher temperatures.

The reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with another important alkyne with the most similar structure to acetylene, propyne $\mathrm{CH}_{3} \mathrm{CCH}$, has also been studied previously. However, compared to the $\mathrm{C}_{2} \mathrm{H}_{2}$ $+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ reaction, the studies are relatively scarce. ${ }^{16-22}$ The total reaction rate constants at room temperature were determined to be at the magnitude of $10^{-13} \mathrm{~cm}^{3}$ molecule ${ }^{-1} \mathrm{~s}^{-1} .{ }^{20-22}$

In an early crossed molecular beam experiment, Kanofsky et al. observed seven reaction channels, including five major ones ( $2 \mathrm{a}-2 \mathrm{e}$ ), for the reaction of $\mathrm{CH}_{3} \mathrm{CCH}$ with $\mathrm{O}\left({ }^{3} \mathrm{P}\right) .{ }^{17}$

$$
\begin{gather*}
\mathrm{CH}_{3} \mathrm{CCH}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{C}_{3} \mathrm{H}_{3} \mathrm{O}+\mathrm{H}  \tag{2a}\\
\mathrm{CH}_{3} \mathrm{CCH}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{C}_{2} \mathrm{HO}+\mathrm{CH}_{3}  \tag{2b}\\
\mathrm{CH}_{3} \mathrm{CCH}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{C}_{2} \mathrm{H}_{4}+\mathrm{CO}  \tag{2c}\\
\mathrm{CH}_{3} \mathrm{CCH}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{C}_{2} \mathrm{H}_{2} \mathrm{O}+\mathrm{CH}_{2}  \tag{2d}\\
\mathrm{CH}_{3} \mathrm{CCH}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{C}_{2} \mathrm{H}_{3}+\mathrm{HCO}  \tag{2e}\\
\mathrm{CH}_{3} \mathrm{CCH}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{CO}+\mathrm{H}_{2}  \tag{2f}\\
\mathrm{CH}_{3} \mathrm{CCH}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{C}_{3} \mathrm{H}_{3}+\mathrm{OH} \tag{2~g}
\end{gather*}
$$

Channels 2 a and 2 b are apparent displacement reactions and are analogous to the channel 1 a in the $\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ reaction, while channels 2 c and 2 f forming CO are analogous to 1 b . It was believed that the first produced ethylidene radical $\mathrm{CH}_{3} \mathrm{CH}$, analogous to the product $\mathrm{CH}_{2}$ in the $\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ reaction, could either rearrange to yield ethylene or decompose to yield acetylene and $\mathrm{H}_{2} \cdot{ }^{16,17,23}$ Another two major channels, the $\mathrm{CH}_{2}$ channel (2d) and the HCO channel (2e), are new channels without analogy to the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{C}_{2} \mathrm{H}_{2}$ reaction. The H -abstraction channel forming OH , channel 2 g , contributed just a small fraction for the total reaction. In order to determine the source of products of $\mathrm{H}, \mathrm{CH}_{3}, \mathrm{HCO}$ as well as OH , deuterated propyne $\mathrm{CH}_{3} \mathrm{CCD}$ and $\mathrm{CD}_{3} \mathrm{CCH}$ were, respectively, reacted with $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ to produce hydrogenated and deuterated products in their experiment.
In another more recent experiment by Bersohn et al. by means of LIF spectroscopy, ${ }^{18} \mathrm{H}, \mathrm{CO}$, and $\mathrm{H}_{2}$ were observed as products from the reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with propyne. The H and CO channels were determined to account for $48 \%$ yield of the total reaction. Moreover, possible mechanism was postulated based on the measurement of the product energy distribution. The CO product was found to be rotationally cold, from which they inferred that the initially formed triplet ketocarbene might cross to a singlet state and isomerize to a substituted ketene by a $1,2-\mathrm{H}$ shift which then dissociated through a linear $\mathrm{C}-\mathrm{C}-\mathrm{O}$ transition state. The $\mathrm{H}_{2}$ with $\mathrm{J} \leq 3$ and rotational temperature of about 400 K was considered to be from the decomposition of $\mathrm{CH}_{3} \mathrm{CH}$ or its isomer $\mathrm{C}_{2} \mathrm{H}_{4}$. In addition, H atoms have a lower translational energy for $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+$ propyne reaction than that for the reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with $\mathrm{C}_{2} \mathrm{H}_{2}$. They thought this is due to the presence of nine additional modes which effectively cool the remainder of the ketocarbene.

Evidently, previous experiments have shown that the product channels for the reaction of propyne with $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ are far more complex than its prototype reaction of acetylene with $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$. The electrophilic $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ is added to the triple bond of alkyne forming a highly energized ketocarbene intermediate. It is expected that this high energy would allow direct decomposition or internal migration of H atoms before dissociation over even large barriers leading to a variety of open channels. The substitution of a methyl group to acetylene introduces additional degrees of freedom for the unimolecular decomposition and internal molecular rearrangements of the reaction intermediates. Thus, rich and complex product channels and reaction mechanisms are anticipated for the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+$ propyne reaction. To our best knowledge, no theoretical calculations have ever been performed for this reaction before.
In this work for the reaction system of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with propyne, the $\operatorname{CCSD}(\mathrm{T}) / / \mathrm{B} 3 \mathrm{LYP} / 6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level of theoretical calcula-
tions are performed for both the triplet and singlet potential energy surfaces and the minimum energy crossing point between the two surfaces are located with the Newton-Lagrange method. On the basis of the theoretical calculations the reaction channels and mechanisms are clarified. Interesting questions are explored as to the site selectivity of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ atom being added to which carbon atom of the triple bond, different types of internal H -atom migrations including $1,2-\mathrm{H}$ shift, $3,2-\mathrm{H}$ shift, and $3,1-\mathrm{H}$ shift involved in the reaction, and the nonadiabatic reaction pathways resulted from the intersystem crossing from triplet state to singlet state.

In addition to the theoretical calculation, experiments detecting the CO products are performed by means of step-scan timeresolved Fourier transform infrared emission spectroscopy (TRFTIR). In the previous experiments by Bersohn et al. ${ }^{18}$ with VUV LIF spectroscopy, CO was found to be almost entirely populated in the $v=0$ vibrational ground state (the $v=1 / v=$ 0 ratio of CO population was 0.05 ). It is obscure why CO was absent of vibrational energy despite the fact that the energy is initially localized on the newly formed $\mathrm{C}-\mathrm{O}$ bond and the proposed pathway to produce $\mathrm{C}_{2} \mathrm{H}_{4}+\mathrm{CO}$ on the singlet surface releases the largest amount of energy of $115.1 \mathrm{kcal} \mathrm{mol}^{-1}$ among all possible channels. In this work, by using TR-FTIR emission spectroscopy, a complete detection of the CO product at each vibrationally excited level is fulfilled, from which the vibrational energy disposal of CO is determined. The experiment shows that some vibrational energy, $6.2 \mathrm{kcal} \mathrm{mol}^{-1}$, is released into the CO vibration with CO being vibrationally excited up to $v$ $=5$. The average vibrational energy is found to consist with the statistical partition of the singlet $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ channel, but not with the triplet $\mathrm{CH}_{3} \mathrm{CH}+\mathrm{CO}$ or $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{3} \mathrm{~A}\right)+\mathrm{CO}$ channel. In combination with the present calculation results, the CO product is identified to arise mainly from the singlet methylketene ( ${ }^{1} \mathrm{CH}_{3} \mathrm{CHCO}$ ) dissociation following the intersystem crossing of the triplet ketocarbene adduct $\left({ }^{3} \mathrm{CH}_{3} \mathrm{CCHO}\right)$, i.e., the $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ channel. Fast intersystem crossing via the minimum energy crossing point of the triplet and singlet surfaces is shown to play significant roles resulting into nonadiabatic pathways for this reaction.

## 2. Computational and Experimental Methods

Computationally, the geometries of the reactants, products, various intermediates, and transition states are optimized using the hybrid density functional theory, i.e., Becke's threeparameter nonlocal exchange functional with the nonlocal correlation functional of Lee, Yang, Parr (B3LYP) with the standard $6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ basis sets. ${ }^{24,25}$ For the current reaction involving eight atoms, the B3LYP/6-311G(d, p) level of theory is a balanced method considering the computational efficiency and accuracy. Harmonic vibrational frequencies and the zeropoint energies (ZPE) are calculated at the same level with the optimized geometries. The intermediates are characterized by all the real frequencies. The transition states are confirmed by only one imaginary frequency. Connections of the transition states between two local minima have been confirmed by intrinsic reaction coordinate (IRC) calculations at the B3LYP/ $6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level. ${ }^{26}$ To obtain more reliable energetic data, single-point electronic energies are calculated at the $\operatorname{CCSD}(\mathrm{T}) /$ $6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level using the B3LYP/6-311G(d, p) optimized geometries. All of the theoretical calculations are performed with the Gaussian 03 program package. ${ }^{27}$ The minimum energy crossing point (MECP) on the intersection seam is located at the B3LYP/6-311G(d, p) level using the Newton-Lagrange method, which was introduced by Koga and Morokuma ${ }^{28}$ to
find the point where the energy is the lowest on the $(f-1)$ dimensional hypersurface of seam between two $f$-dimensional potential energy surfaces. A homemade program is used for this purpose. This method has been applied successfully in locating the MECP for the reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with isobutene. ${ }^{43}$

Experimentally, the reaction products are monitored by stepscan, time-resolved Fourier transform emission (TR-FTIR) spectroscopy which is an effective technique probing IR-active reaction products in real time due to its multiplex advantage and nanosecond time resolution. ${ }^{29}$ Step-scan FTIR spectrometer is commercially available but requires significant modification for coupling with pulsed laser to implement time-resolve IR spectral measurements. The details of the TR-FTIR spectrometer and the reaction chamber have been described previously. ${ }^{30,31}$ Briefly, the instrument comprises a Nicolet Nexus 870 stepscan FTIR spectrometer, Spectra Physics (Quanta- Ray PRO 230) Nd:YAG laser, and a pulse generator (Stanford Research DG535) to initiate the laser pulse and achieve synchronization of the laser with data collection, two digitizers (internal 100 KHz 16-bit digitizer and external 100 MHz 14-bit GAGE 8012A digitizer) which offer fast time resolution and a wide dynamic range as needed, and a personal computer to control the whole experiment. The detector used in this experiment is a liquid nitrogen cooled InSb detector.
The reaction is initiated in a stainless steel flow reaction chamber with a 20 -L-volume. $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ atoms were produced via the laser photodissociation of $\mathrm{NO}_{2}$ at 355 nm (Nd:YAG laser, Spectra Physics Quanta-Ray PRO 230, pulse energy $\sim 100 \mathrm{~mJ}$ ). Only the ground-state $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ atoms are produced following the 355 nm photolysis of $\mathrm{NO}_{2}$ and there is no interference from $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ atoms. Typically 200 mTorr of propynes $(\geq 99.5 \%)$ and 150 mTorr of $\mathrm{NO}_{2}(\geq 99.9 \%)$ without buffer gas enter the flow chamber 1 cm above the photolysis beam via needle valves. The chamber is pumped by an $8 \mathrm{~L} \mathrm{~s}^{-1}$ mechanical pump and the stagnation pressure of the chamber is measured by a MKS capacitance monometer. With the photolysis volume in the cylinder of 250 mL , it is expected that the flow rate is fast enough to replenish the sample at each laser pulse running normally at a repetition rate of $10 \mathrm{~Hz}^{30,31}$ The constant pressure of sample is maintained by adjusting the pumping speed and the needle valves. Transient infrared emission is collected by a pair of gold-coated White-Cell spherical mirrors and collimated by a $\mathrm{CaF}_{2}$ lens to the step-scan FTIR spectrometer.

## 3. Results and Discussion

3.1. Theoretical Calculations of the Product Channels and Mechanisms. The optimized geometries of various reactants, intermediates, transition states and products at the B3LYP/6$311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level are shown in Figure 1. The vibrational modes and imaginary frequencies of the transition states are also indicated. The zero-point vibrational energies at the B3LYP/ $6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level, the single-point energies computed with the $\operatorname{CCSD}(\mathrm{T})$ method, and the relative energies by taking the energy of the reactants as zero are listed in Table 1. The energy difference between the reactants and products computed at the $\operatorname{CCSD}(\mathrm{T}) / / \mathrm{B} 3 \mathrm{LYP} / 6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level agrees well with the available experimental reaction enthalpies, demonstrating that the present level of calculation can provide reliable energetic and mechanistic information for the reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with propyne.
3.1.1. The Minimum Energy Crossing Point between the Singlet and Triplet States. The reaction starts on a triplet surface with the addition of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ to the triple bond of propyne forming a ketocarbene $\mathrm{CH}_{3} \mathrm{CCHO}$ * from which a variety of reaction
channels are open. Also, previous experiments ${ }^{17}$ indicated the existence of nonadiabatic pathways forming $\mathrm{C}_{2} \mathrm{H}_{4}+\mathrm{CO}$ after the initial triplet ketocarbene crosses to a singlet state. Therefore, reaction pathways on both triplet and singlet potential energy surfaces as well as the minimum energy crossing point (MECP) between the two states are calculated.

It has been shown that the characterization of the $\mathrm{MECP}^{35-41}$ plays an important role in the investigation of the chemical reaction mechanism. For polyatomic molecules, there may be many intersections between two potential energy surfaces. The MECP on the intersection seam is very important and is usually considered as a "transition state" for the nonadiabatic process.

Koga and Morokuma ${ }^{28}$ introduced the Newton-Lagrange method for the search of the MECP, which has the same geometry and energy for the singlet and triplet states. The energies, energy gradients, and Hessian matrixes of both singlet and triplet states are calculated, and the lowest-energy point is found on the seam of intersection at the B3LYP/6-311G(d,p) level. The obtained geometry for the MECP is between the two equilibrium geometries of ${ }^{3}$ IM1 and ${ }^{1}$ IM9 as shown in Figure 2. The energies and energy gradients of the MECP in the triplet and singlet states are listed in Table 2. The energy gradients of the MECP are not zero, unlike the optimization result of an equilibrium or a transition state geometry. The energy gradients of the MECP in the singlet state are proportional to that in the triplet state, and the ratio equals $-\lambda /(1-\lambda)$, where $\lambda$ is the Lagrange multiplier. These characteristics are shown to be reasonable, which is a good check for the obtained MECP.

The chemical reaction pathway can be found by the intrinsic reaction coordinate (IRC) method, ${ }^{26}$ which is used to search the reaction paths of the intersystem crossing through the intersection. Figure 3 shows the minimum energy path from ${ }^{3}$ IM1 to ${ }^{1}$ IM 9 through the MECP varying with the $2 \mathrm{C}-1 \mathrm{C}$ bond distance and $3 \mathrm{C}-1 \mathrm{C}-2 \mathrm{C}$ bond angle at the B3LYP/6-311G(d,p) level. The other coordinates are optimized. It can be seen that the MECP connects ${ }^{3} \mathrm{IM} 1$ and ${ }^{1} \mathrm{IM} 9$ properly. There is no experimental information about the energy barrier between ${ }^{3} \mathrm{IM} 1$ and the MECP. The calculations of this work show that this barrier is $13.9 \mathrm{kcal} \mathrm{mol}^{-1}$ above ${ }^{3} \mathrm{IM} 1$ at the $\operatorname{CCSD}(\mathrm{T}) / 6-$ $311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level. We will show later that the MECP plays a key role and the intersystem crossing has the lowest barrier among all reaction routes in the triplet state.
3.1.2. Adiabatic Pathways on the Triplet Potential Energy Surface. The adiabatic reaction pathways occurring on the triplet surface are explored and shown in Figure 4. The reaction starts with the addition of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ to either of the carbon atoms of the triple bond (the terminal carbon atom is referred to the C 1 atom and the other is the C 2 atom) or abstraction of an H atom from the methyl group. These three pathways are discussed separately and their relative importance is indicated based on the barrier height of the transition state (TS) as well as the rigidity of the TS, i.e., its "looser" or "tighter" entropic characteristic.

O-Atom Addition to the C1 Atom. Addition of the O atom to the C 1 atom can occur via TS1 leading to the initial ketocarbene adduct IM1 $\left(\mathrm{CH}_{3} \mathrm{CCHO}\right)$ lying $48.2 \mathrm{kcal} \mathrm{mol}{ }^{-1}$ below the reactants. This step faces a barrier height of $5.4 \mathrm{kcal} \mathrm{mol}^{-1}$. Starting from IM1, there are four different reaction pathways as follows: (i) breaking the $\mathrm{C} 1-\mathrm{H}$ bond to yield $\mathrm{CH}_{3} \mathrm{CCO}+\mathrm{H}$ atom via TS2 with a barrier of $38.9 \mathrm{kcal} \mathrm{mol}^{-1}$; (ii) losing a methyl H atom via TS3 to form $\mathrm{CH}_{2} \mathrm{CCHO}$ plus H atom, facing a barrier of $47.1 \mathrm{kcal} \mathrm{mol}^{-1}$; (iii) isomerizing to IM2 $\left(\mathrm{CH}_{3} \mathrm{CHCO}\right)$ by a $1,2-\mathrm{H}$ shift via TS 4 with a barrier of 44.5 $\mathrm{kcal} \mathrm{mol}^{-1}$; and (iv) rearranging to IM3 $\left(\mathrm{CH}_{2} \mathrm{CHCHO}\right)$ via TS5 through a $3,2-\mathrm{H}$ shift with a barrier of $46.7 \mathrm{kcal} \mathrm{mol}^{-1}$. The



$\mathrm{CH}_{3} \mathrm{CCO}+\mathrm{H}$
$\mathrm{C}-1.127 \mathrm{O}$



$$
\mathrm{CH}_{2} \mathrm{CO}+\mathrm{CH}_{2}
$$


$\mathrm{CH}_{2} \mathrm{CH}+\mathrm{HCO}$

$\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{CO}+\mathrm{H}_{2}$


IM4


IM8


IM9


TS1, Freq $=268.5 \mathrm{i}$


IM10


TS 2, Freq $=493.3 \mathrm{i}$


IM11


TS3, Freq $=367.1 \mathrm{i}$

Figure 1. Part 1 of 2.


TS4, Freq $=1701.6 \mathrm{i}$
TS5, Freq $=1779.0 \mathrm{i}$


TS9, Freq $=1649.5 \mathrm{i}$


TS13, Freq $=370.1 \mathrm{i}$


TS16, Freq $=418.1 i$
TS17, Freq $=2008.7 \mathrm{i}$


TS21, Freq $=229.3 \mathrm{i}$

TS24, Freq $=1449.4 \mathrm{i}$



TS25, Freq $=1910.7 \mathrm{i}$


TS6, Freq $=139.3 i$


TS10, Freq $=298.6 \mathrm{i}$


TS14, Freq $=455.1 \mathrm{i}$


TS18, Freq $=298.3 \mathrm{i}$


TS22, Freq $=977.1 \mathrm{i}$


TS26, Freq $=871.5 \mathrm{i}$


TS7, Freq $=261.9 \mathrm{i}$


TS11, Freq $=366.9 \mathrm{i}$


TS15, Freq $=2105.6 \mathrm{i}$


TS19, Freq $=2125.3 \mathrm{i}$


TS23, Freq $=1005.8 \mathrm{i}$


TS27, Freq $=231.8 \mathrm{i}$

Figure 1. Part 2 of 2 . Optimized geometries of reactants, intermediates, transition states, and products for the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{CH}_{3} \mathrm{CCH}$ reaction at the $\mathrm{B} 3 \mathrm{LYP} /$ $6-311 G(d, p)$ level. The vibrational modes and the imaginary frequencies of the transition states are also indicated. Bond lengths are in angstrom, bond angles are in degrees and imaginary frequencies are in $\mathrm{cm}^{-1}$.

TABLE 1: Zero-Point Energies (ZPE, hartree), Total Energies (TE, hartree) and Relative Energies (RE, kcal mol ${ }^{-1}$ ) Obtained by Taking the Energy of Reactants As Zero for Various Species in the $\mathbf{O}\left({ }^{3} \mathbf{P}\right)+\mathbf{C H}_{3} \mathbf{C C H}$ Reaction

| species | ZPE ${ }^{\text {a }}$ | TE ${ }^{b}$ | RE | $\Delta H_{298 \mathrm{~K}}{ }^{\text {c }}$ | species | $\mathrm{ZPE}^{a}$ | TE ${ }^{b}$ | RE |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{C}_{3} \mathrm{H}_{4}+\mathrm{O}$ |  | -191.2371891 | 0.0 |  | TS2( ${ }^{3} \mathrm{~A}$ ) | 0.049495 | -191.2519642 | -9.3 |
| $\mathrm{CH}_{3} \mathrm{CCO}+\mathrm{H}$ |  | -191.2585157 | -13.4 |  | TS3 ( ${ }^{3} \mathrm{~A}$ ) | 0.047668 | -191.2389389 | -1.1 |
| $\mathrm{CH}_{2} \mathrm{CCHO}+\mathrm{H}$ |  | -191.2437677 | -4.1 |  | TS4( ${ }^{3} \mathrm{~A}$ ) | 0.052819 | -191.2430541 | -3.7 |
| $\mathrm{CH}_{2} \mathrm{CHCO}+\mathrm{H}$ |  | -191.2740194 | -23.1 |  | TS5 ( ${ }^{3} \mathrm{~A}$ ) | 0.053122 | -191.2395944 | -1.5 |
| $\mathrm{CH}_{3} \mathrm{CH}+\mathrm{CO}$ |  | -191.3134707 | -47.9 |  | TS6( ${ }^{3} \mathrm{~A}$ ) | 0.052978 | -191.3084379 | -44.7 |
| $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{3} \mathrm{~A}\right)+\mathrm{CO}$ |  | -191.3200920 | -52.0 |  | TS7( ${ }^{3} \mathrm{~A}$ ) | 0.049455 | -191.2696828 | -20.4 |
| $\mathrm{CH}_{3}+\mathrm{HCCO}$ |  | -191.2724856 | -22.1 | -26.4 | TS8( ${ }^{3} \mathrm{~A}$ ) | 0.051801 | -191.2609387 | -14.9 |
| $\mathrm{CH}_{2} \mathrm{CO}+\mathrm{CH}_{2}$ |  | -191.2683560 | -19.6 | -22.1 | TS9 ( ${ }^{3} \mathrm{~A}$ ) | 0.054685 | -191.2722334 | -22.0 |
| $\mathrm{CH}_{2} \mathrm{CH}+\mathrm{HCO}$ |  | -191.2674938 | -19.0 | -21.9 | TS10( ${ }^{3} \mathrm{~A}$ ) | 0.053108 | -191.3069492 | -43.8 |
| $\mathrm{CH}_{2} \mathrm{CCH}+\mathrm{OH}$ |  | -191.2442603 | -4.4 | -11.8 | TS11( ${ }^{3} \mathrm{~A}$ ) | 0.051614 | -191.2547985 | -11.1 |
| $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ |  | -191.4206778 | $-115.1$ | -117.9 | TS12( ${ }^{3} \mathrm{~A}$ ) | 0.052722 | -191.2587638 | -13.5 |
| $\mathrm{H}_{2} \mathrm{CCCO}+\mathrm{H}_{2}$ |  | -191.3431980 | -66.5 |  | TS13( ${ }^{3} \mathrm{~A}$ ) | 0.055756 | -191.2264549 | 6.7 |
| $\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{CO}+\mathrm{H}_{2}$ |  | -191.3594572 | -76.7 | $-76.0$ | TS14( ${ }^{3} \mathrm{~A}$ ) | 0.052590 | -191.2584734 | -13.4 |
| $\mathrm{IM} 1\left({ }^{3} \mathrm{~A}\right)$ | 0.058492 | -191.3139894 | -48.2 |  | TS15( ${ }^{3} \mathrm{~A}$ ) | 0.051747 | -191.2501285 | -8.1 |
| IM2( ${ }^{3} \mathrm{~A}$ ) | 0.058962 | -191.3400544 | -64.5 |  | TS16( ${ }^{3} \mathrm{~A}$ ) | 0.051536 | -191.2528385 | -9.8 |
| IM3( ${ }^{3} \mathrm{~A}$ ) | 0.056508 | -191.3174641 | -50.4 |  | TS17( ${ }^{3} \mathrm{~A}$ ) | 0.052029 | -191.2151272 | 13.8 |
| IM4( ${ }^{3} \mathrm{~A}$ ) | 0.056902 | -191.3258612 | -55.6 |  | TS18( ${ }^{3} \mathrm{~A}$ ) | 0.056526 | -191.2910346 | -33.8 |
| IM5 ( ${ }^{3} \mathrm{~A}$ ) | 0.057878 | -191.3140792 | -48.2 |  | TS19( ${ }^{3} \mathrm{~A}$ ) | 0.050946 | -191.2354739 | 1.1 |
| IM6( ${ }^{3} \mathrm{~A}$ ) | 0.056900 | -191.3317251 | -59.3 |  | $\mathrm{TS} 20\left({ }^{3} \mathrm{~A}\right)$ | 0.049691 | -191.2186171 | 11.7 |
| IM7( ${ }^{3} \mathrm{~A}$ ) | 0.057198 | -191.2936118 | -35.4 |  | TS21( ${ }^{1} \mathrm{~A}$ ) | 0.056008 | -191.3005295 | -39.7 |
| $\mathrm{IM} 8\left({ }^{3} \mathrm{~A}\right)$ | 0.057000 | -191.2939019 | -35.6 |  | TS22( ${ }^{1} \mathrm{~A}$ ) | 0.055347 | -191.2940147 | -35.7 |
| $\operatorname{IM} 9\left({ }^{1} \mathrm{~A}\right)$ | 0.057599 | -191.3014430 | -40.3 |  | TS23( ${ }^{1} \mathrm{~A}$ ) | 0.055468 | -191.2966087 | -37.3 |
| $\mathrm{IM} 10\left({ }^{1} \mathrm{~A}\right)$ | 0.060861 | -191.4126978 | -110.1 |  | TS24( ${ }^{1} \mathrm{~A}$ ) | 0.055173 | -191.2906739 | -33.6 |
| $\mathrm{IM} 11\left({ }^{1} \mathrm{~A}\right)$ | 0.061178 | -191.4094665 | -108.1 |  | TS25 ( ${ }^{1} \mathrm{~A}$ ) | 0.049446 | -191.2670007 | -18.7 |
| $\mathrm{IM} 12\left({ }^{1} \mathrm{~A}\right)$ | 0.061094 | -191.4117219 | -109.5 |  | TS26( ${ }^{1} \mathrm{~A}$ ) | 0.048491 | -191.2813782 | -27.7 |
| TS1 $\left({ }^{3} \mathrm{~A}\right)$ | 0.055729 | -191.2285956 | 5.4 |  | TS27( ${ }^{1} \mathrm{~A}$ ) | 0.060042 | -191.4009522 | -102.8 |

[^1]

Figure 2. The optimized geometries and the numbering of atoms for the minimum energy crossing point (MECP) at the B3LYP/6-311G(d,p) level.
two H -atom displacement pathways, (i) and (ii), are more favorable than the two H -atom migration pathways, (iii) and (iv), due to the looser entropic character of the transition states.

The methylketene intermediate, IM2 $\left(\mathrm{CH}_{3} \mathrm{CHCO}\right)$, lying 64.5 $\mathrm{kcal} \mathrm{mol}^{-1}$ below the initial reactants, can decompose to triplet $\mathrm{CH}_{3} \mathrm{CH}$ plus CO via TS6 by $\mathrm{C} 1-\mathrm{C} 2$ bond rupture, or to $\mathrm{CH}_{2} \mathrm{CHCO}+\mathrm{H}$ via TS7 by losing a H atom, or to $\mathrm{CH}_{3}$ plus HCCO via TS8 by C2-C3 bond cleavage. These pathways confront barriers of $19.8 \mathrm{kcal} \mathrm{mol}^{-1}, 44.1 \mathrm{kcal} \mathrm{mol}^{-1}$ and 49.6 $\mathrm{kcal} \mathrm{mol}{ }^{-1}$, respectively. In addition, IM2 can also isomerize to IM4 $\left(\mathrm{CH}_{2} \mathrm{CH}_{2} \mathrm{CO}\right)$ via TS9 with a barrier of $42.5 \mathrm{kcal} \mathrm{mol}^{-1}$. Then IM4, $55.6 \mathrm{kcal} \mathrm{mol}^{-1}$ below the initial reactants, can either break $\mathrm{C} 1-\mathrm{C} 2$ bond to yield triplet $\mathrm{C}_{2} \mathrm{H}_{4}$ plus CO via TS10 overcoming a barrier of $11.8 \mathrm{kcal} \mathrm{mol}^{-1}$, or decompose to ketene plus $\mathrm{CH}_{2}$ via TS11 by $\mathrm{C} 3-\mathrm{C} 2$ bond rupture, facing a barrier of $44.5 \mathrm{kcal} \mathrm{mol}^{-1}$. Among these four pathways starting from IM2, the decomposition channel to $\mathrm{CH}_{3} \mathrm{CH}$ plus CO should be dominant, considering its lowest barrier height.

The aldehyde intermediate, IM3 ( $\mathrm{CH}_{2} \mathrm{CHCHO}$ ), 50.4 kcal $\mathrm{mol}^{-1}$ below the initial reactants, can decompose to $\mathrm{CH}_{2} \mathrm{CH}$ plus HCO via TS12 over a barrier of $36.9 \mathrm{kcal} \mathrm{mol}^{-1}$.

O-Atom Addition to the C2 Atom. The addition of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ to C 2 atom of the triple bond leads to another diradical adduct, IM5 $\left(\mathrm{CH}_{3} \mathrm{COCH}\right)$, via TS13 after surmounting a barrier of 6.7
kcal $\mathrm{mol}^{-1}$. This barrier is just a little higher than that of O-addition to C 1 atom, indicating that the addition to C 2 can somewhat compete with the addition to C 1 atom. The identical binding energy of IM5 to IM1, $48.2 \mathrm{kcal} \mathrm{mol}{ }^{-1}$, also suggests that no preference is expected for the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ addition to either C 1 or C 2 atom. It can be concluded that there is no strong site selectivity when $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ is added to the triple bond carbon atoms considering the nearly identical barrier height and binding energies of the adducts. Similar behavior of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ has also been observed experimentally ${ }^{42}$ and verified later by theoretical calculations ${ }^{43}$ when studying the reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with alkenes.

Starting from IM5 $\left(\mathrm{CH}_{3} \mathrm{COCH}\right)$, there are three pathways, namely: (i) direct decomposition to $\mathrm{CH}_{3}$ plus HCCO via TS14 with a barrier of $34.8 \mathrm{kcal} \mathrm{mol}^{-1}$; (ii) isomerization to IM6 $\left(\mathrm{CH}_{2} \mathrm{COCH}_{2}\right)$ by 3,1-H shift via TS15 overcoming a barrier of $40.1 \mathrm{kcal} \mathrm{mol}^{-1}$, and then subsequent decomposition to ketene $\mathrm{CH}_{2} \mathrm{CO}$ plus $\mathrm{CH}_{2}$ via TS16 with a barrier of $49.5 \mathrm{kcal} \mathrm{mol}^{-1}$; (iii) rearrangement to IM7 via TS17 with the subsequent isomerization and decomposition leading finally to $\mathrm{CH}_{2} \mathrm{CO}+$ $\mathrm{CH}_{2}$. Because TS17 lies $13.8 \mathrm{kcal} \mathrm{mol}^{-1}$ above the reactants, this reaction route is energetically unfavorable. No doubt, the $\mathrm{CH}_{3}+\mathrm{HCCO}$ channel is the most favorable of all the three pathways starting from the adduct IM5 $\left(\mathrm{CH}_{3} \mathrm{COCH}\right)$. This can be explained by the fact that when oxygen atom being added to C 2 atom, the weakest $\mathrm{C} 3-\mathrm{C} 2$ bond is subject to cleavage the most easily.

The H-Abstraction. The $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ atom can abstract a methyl H atom in propyne via TS20 yielding hydroxyl OH plus propargyl radical $\mathrm{CH}_{2} \mathrm{CCH}$. Compared to the entrance barriers of O addition reactions, TS20 is 6.3 and $5.0 \mathrm{kcal} \mathrm{mol}^{-1}$ higher in energy than TS1 and TS13, respectively. Consequently, the H -abstraction channel cannot compete with the O -addition channels. An alternate acetylentic H abstraction channel has been

TABLE 2: Energies (in Hartree) and Energy Gradients (g, in hartree/bohr) of the Minimum Energy Crossing Point (MECP) in the Triplet (T) and Singlet (S) States at the B3LYP/6-311G(d,p) Level ${ }^{a}$

${ }^{a}$ Bond distances $\left(B_{\mathrm{n}}\right)$, angles $\left(A_{\mathrm{n}}\right)$, and dihedrals $\left(D_{\mathrm{n}}\right)$ are defined according to the numbering of the MECP in Figure 2.

$2 \mathrm{C}-1 \mathrm{C}$
Figure 3. Minimum energy path from ${ }^{3} \mathrm{IM} 1$ to ${ }^{1} \mathrm{IM} 9$ through the minimum energy crossing point (MECP) varying with $2 \mathrm{C}-1 \mathrm{C}$ bond distance and $3 \mathrm{C}-1 \mathrm{C}-2 \mathrm{C}$ bond angle at the B3LYP/6-311G(d,p) level. (Other coordinates are optimized.)
ruled out since no transition state can be located for this type of H abstraction reaction.

Overall, the reaction on the triplet potential energy surface can be summarized by two types of reactions, i.e., Habstraction and O -addition. The H -abstraction channels should contribute only a small fraction for the total reaction
due to the high energy barrier ( $11.7 \mathrm{kcal} \mathrm{mol}^{-1}$ ) it has to surmount. In contrast, the O -addition channels should account for most of the total reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with propyne and the addition can occur on both of the carbon atoms of the triple bond leading first to the diradical adducts IM1 $\left(\mathrm{CH}_{3} \mathrm{CCHO}\right)$ and IM5 $\left(\mathrm{CH}_{3} \mathrm{COCH}\right)$. Subsequently, the energized adducts IM1 and IM5 are subject to direct decomposition or H migrations (including $1,2-\mathrm{H}$ shift, $3,2-\mathrm{H}$ shift and $3,1-\mathrm{H}$ shift) forming various isomeric intermediates IM2 $\left(\mathrm{CH}_{3} \mathrm{CHCO}\right)$, IM3 $\left(\mathrm{CH}_{2} \mathrm{CHCHO}\right)$, and IM6 $\left(\mathrm{CH}_{2} \mathrm{COCH}_{2}\right)$, which decompose further to final products. Compared to the direct decomposition of the diradical adducts forming radical pairs $\mathrm{CH}_{3} \mathrm{CCO}+\mathrm{H}, \mathrm{CH}_{2} \mathrm{CCHO}+\mathrm{H}$ and $\mathrm{CH}_{3}+\mathrm{HCCO}$, the H migration channels forming eventually $\mathrm{CH}_{2} \mathrm{CH}+\mathrm{HCO}$, $\mathrm{CH}_{2} \mathrm{CO}+\mathrm{CH}_{2}, \mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{3} \mathrm{~A}\right)+\mathrm{CO}$ and $\mathrm{CH}_{3} \mathrm{CH}+\mathrm{CO}$ are less competitive because the H migration requires surmounting higher barriers and more steps of molecular rearrangements. Also, the H and $\mathrm{CH}_{3}$ displacement channels are further facilitated due to the looser entropic character of the direct decomposition. The most favorable reaction channels on the triplet surface are therefore concluded to be $\mathrm{CH}_{3} \mathrm{CCO}+\mathrm{H}$ $\left(\mathrm{CH}_{2} \mathrm{CCHO}+\mathrm{H}\right)$ and $\mathrm{CH}_{3}+\mathrm{HCCO}$ occurring via the direct decomposition of the diradical adducts IM1 and IM5. Interestingly, three types of H migrations are found to play equally important roles following the initial O -addition if considering their nearly identical barrier heights, i.e., 44.5 $\mathrm{kcal} \mathrm{mol}^{-1}$ for the $1,2-\mathrm{H}$ shift from IM1 to IM2, 46.7 kcal


Figure 4. Potential energy profiles of the adiabatic reaction pathways on the triplet surface at the $\operatorname{CCSD}(\mathrm{T}) / / \mathrm{B} 3 \mathrm{LYP} / 6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level. Key: (a) reaction pathways starting from the O -atom addition to the C 1 atom of the triple bond; (b) reaction pathways starting from the O -atom addition to the C 2 atom of the triple bond; as well as the H -abstraction channel.
$\mathrm{mol}^{-1}$ for the $3,2-\mathrm{H}$ shift from IM1 to IM3, and 40.1 kcal $\mathrm{mol}^{-1}$ for the $3,1-\mathrm{H}$ shift from IM5 to IM6.
3.1.3. Nonadiabatic Pathways on the Singlet Potential Energy Surface. As discussed in section 3.1.1, the minimum energy crossing point between the triplet and singlet states has been located to be $13.9 \mathrm{kcal} \mathrm{mol}^{-1}$ above ${ }^{3} \mathrm{IM} 1$, and thus the intersystem crossing has the lowest energy barrier among the reaction routes of ${ }^{3} \mathrm{IM} 1$. Through MECP, the triplet ketocarbene ${ }^{3}$ IM1 crosses into its singlet state ${ }^{1}$ IM9 efficiently, from which a variety of nonadiabatic reaction pathways are open as shown in Figure 5.

On the singlet surface through a nearly barrierless $1,2-\mathrm{H}$ shift process, ${ }^{1}$ IM9 rearranges to the highly energized singlet methylketene, IM10, from which three channels are open: (i) cleaving the $\mathrm{C} 1-\mathrm{C} 2$ bond with the concerted $3,2-\mathrm{H}$ shift via TS24 to yield $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$; (ii) rearranging to IM 12 which undergoes subsequently $\mathrm{H}_{2}$ elimination while breaking the $\mathrm{C} 1-\mathrm{C} 2$ bond to yield ultimately $\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{H}_{2}+\mathrm{CO}$; (iii) subsequent rearrangement of IM12 to IM11 which undergoes $\mathrm{H}_{2}$ elimination leading to $\mathrm{H}_{2}+\mathrm{H}_{2} \mathrm{CCCO}$. IM11 can also be formed via the $3,2-\mathrm{H}$ shift of IM9 with a low barrier of $4.6 \mathrm{kcal} \mathrm{mol}^{-1}$. The rate-limiting steps of the three channels, TS24 corresponding to the $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ channel, is the lowest in energy barrier
compared to TS26 (corresponding to $\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{H}_{2}+\mathrm{CO}$ channel) and TS25 (corresponding to $\mathrm{H}_{2}+\mathrm{H}_{2} \mathrm{CCCO}$ channel). Therefore, the $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ channel dominates on the singlet surface and the $\mathrm{H}_{2}+\mathrm{H}_{2} \mathrm{CCCO}$ and $\mathrm{C}_{2} \mathrm{H}_{2}+\mathrm{H}_{2}+\mathrm{CO}$ are minor channels.

The nonadiabatic pathway of $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ via the singlet surface are far more exothermic $\left(\Delta H=-115.1 \mathrm{kcal} \mathrm{mol}^{-1}\right)$ than that of the $\mathrm{CH}_{3} \mathrm{CH}+\mathrm{CO}\left(\Delta H=-47.9 \mathrm{kcal} \mathrm{mol}^{-1}\right)$ and $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{3} \mathrm{~A}\right)+\mathrm{CO}\left(\Delta H=-52.0 \mathrm{kcal} \mathrm{mol}^{-1}\right)$ channels occurring on the triplet surface. In addition, the intersystem crossing through MECP corresponds to the lowest energy barrier among all the pathways of the ketocarbene ${ }^{3} \mathrm{IM} 1$ on the triplet surface. Therefore, CO is most likely produced from the singlet surface judging from the present theoretical calculation.
3.2. Experimental Investigation of the CO Channel. To assist identifying the source of CO produced whether from the triplet or singlet reaction pathways, the CO channel is probed by means of TR-FTIR emission spectra in the present experiment. This technique allows a complete detection of the product CO at each vibrationally excited level because it records the infrared fluorescence due to $v \rightarrow v-1$ vibrational transitions of a whole set of vibrationally excited levels. Further spectral analysis can reveal an almost complete population distribution


Figure 5. Potential energy profiles of the nonadiabatic reaction pathways on the singlet surface at the $\operatorname{CCSD}(\mathrm{T}) / / \mathrm{B} 3 \mathrm{LYP} / 6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ level.


Figure 6. Product TR-FTIR emission spectra from the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{CH}_{3} \mathrm{CCH}$ reaction taken at typical delay times from 6 to $400 \mu$ s after initiation of the reaction by 355 nm laser photolysis. The spectra were collected with a resolution of $0.5 \mathrm{~cm}^{-1}$.
of vibrationally excited levels from which the product vibrational energy disposal can be determined.

In the first reference experiment, when pure propyne was irradiated by 355 nm laser, no IR emission was observed because propynes can not be photodissociated by a 355 nm laser. In the second reference experiment when pure $\mathrm{NO}_{2}$ was irradiated by 355 nm laser, only the IR emission spanning from 1790 to 1970 $\mathrm{cm}^{-1}$ was observed due to the vibrationally excited photofragments of NO.

$$
\mathrm{NO}_{2}+h v(355 \mathrm{~nm}) \rightarrow \mathrm{NO}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)
$$

Fortunately, the NO spectral region is isolated and will not interfere with the detection of the products of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with propynes with the spectral region higher than $1970 \mathrm{~cm}^{-1}$.

When the flowing gaseous mixture of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ precursor, $\mathrm{NO}_{2}$, with propynes was irradiated by 355 nm laser, one rotationally resolved IR emission band spanning from 2000 to $2200 \mathrm{~cm}^{-1}$ was observed as shown in Figure 6. Judging from its spectral position and rotational structure, this band is assigned to the vibrationally excited CO products from the reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with propynes. The IR emission intensity of CO appears in a few microseconds and reaches its maximum intensity approximately at $50 \mu \mathrm{~s}$.

A spectral fitting has been performed for the rotationally resolved CO emission bands using the spectral constants of CO and a nonlinear least-squares fitting program which has been described in detail elsewhere. ${ }^{44}$ Representative fitting result for the $50 \mu$ s spectrum shown in Figure 7 a, demonstrates further


Figure 7. (a) Representative spectral fitting results for the IR emission bands of the product CO. (b) Representative Boltzmann plot of the vibrational distribution of CO. The straight line is the fit of the data to a Boltzmann distribution with a temperature of $3151 \pm 223 \mathrm{~K}$.
that this band is ascribed to CO. The best-fitted rotational temperature is nearly 300 K , the room temperature. This is reasonable because about 230 collisions take place within 50 $\mu \mathrm{s}$ at the total pressure of 350 mTorr with the rotational excitation of the products being quenched completely. To ensure adequate signal-to-noise ratios for detection purposes, the species partial pressures are necessarily relatively high, with which the rotational energy distribution is not possibly obtained because of collisions. The best fitted vibrational populations of CO are 1/0.32/0.09/0.04/0.02 for the vibrational levels of $v=1-5$. The IR emission band of CO consists of five $v \rightarrow v-1$ progressions of rovibrational transitions. As shown in Figure 7 b, the vibrational population can be fitted nicely by a Boltzmann distribution with a vibrational temperature of $T_{\text {vib }}=3151 \pm$ 223 K , corresponding to an average vibrational energy of 6.2 $\mathrm{kcal} \mathrm{mol}^{-1}$. This is presumably close to the nascent vibrational energy because of two reasons. First, with a reaction rate constant at the magnitude of $10^{-13} \mathrm{~cm}^{3}$ molecule ${ }^{-1} \mathrm{~s}^{-1},{ }^{20-22}$ the products are still newly forming at $50 \mu \mathrm{~s}$. Second, the vibrational distribution of CO is not expected to be altered much because the vibrational relaxation of CO is highly inefficient. ${ }^{31}$ It shows here that some vibrational energy is released into the CO vibration. When probing CO with VUV LIF spectroscopy, Bersohn et al. ${ }^{18}$ found that CO was populated almost entirely at $v=0$ level and hence CO is lack of vibrational energy. This is somehow in conflict with the fact that the energy is initially localized on the newly formed $\mathrm{C}-\mathrm{O}$ bond and the large exothermicity of $115.1 \mathrm{kcal} \mathrm{mol}^{-1}$ of the $\mathrm{C}_{2} \mathrm{H}_{4}+\mathrm{CO}$ channel. As an absorption-based spectroscopy technique, LIF is probably not as sensitive as IR emission spectroscopy toward detecting molecules at vibrationally excited levels in some cases. This might be the reason why Bersohn et al. did not observe much CO at $v \geq 1$ levels.

The present theoretical calculation shows that there are two possible pathways producing CO either on the triplet or singlet surface:

$$
\begin{gathered}
\mathrm{CH}_{3} \mathrm{CCH}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{CH}_{3} \mathrm{CH}+\mathrm{CO} \text { or } \mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{3} \mathrm{~A}\right)+\mathrm{CO} \\
\Delta H=-47.9(\text { or }-52.0) \mathrm{kcal} \mathrm{~mol}{ }^{-1} \\
\rightarrow \mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}
\end{gathered}
$$

$$
\Delta H=-115.1 \mathrm{kcal} \mathrm{~mol}^{-1}
$$

Considering the simplest possible statistical partition with the energy democratically released into the 13 vibrational and 3 relative translational modes of the $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ products on the singlet surface, the vibrational energy partitioned into CO should be $7.2 \mathrm{kcal} \mathrm{mol}^{-1}$, while the triplet channels $\mathrm{CH}_{3} \mathrm{CH}$
+CO and $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{3} \mathrm{~A}\right)+\mathrm{CO}$ should yield CO with much less energy ( 3.0 or $3.2 \mathrm{kcal} \mathrm{mol}^{-1}$ ). The average vibrational energy of $6.2 \mathrm{kcal} \mathrm{mol}^{-1}$ extracted from the CO IR emission spectra fits with the statistically predicted energy release from the singlet channel, but not with the triplet channels. This implies that CO should primarily arise from the singlet mechanism.

Interestingly, Bersohn et al. ${ }^{18}$ observed that CO was also extremely cold in rotation and hence they speculated that CO was most likely produced from the decomposition of the singlet methylketene through a collinear $\mathrm{C}-\mathrm{C}-\mathrm{O}$ transition state to avoid torques on CO. Our calculation shows that the linear singlet methylketene (IM10) dissociates to form $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+$ CO through a transition state (TS24) with a bent $\mathrm{C}-\mathrm{C}-\mathrm{O}$ angle of $141^{\circ}$. Although the calculation shows that the transition state leading to CO is not collinear as postulated, ${ }^{18}$ the rotational excitation of CO should not to be large if produced from the slightly bent structure of TS24. Together with the CO vibrational energy disposal determined by the present experiment, it can now be concluded that CO arises mainly from the singlet methylketene ( $\left.{ }^{1} \mathrm{CH}_{3} \mathrm{CHCO}\right)$ dissociation following the intersystem crossing of the triplet ketocarbene adduct $\left({ }^{3} \mathrm{CH}_{3} \mathrm{CCHO}\right)$. This is reasonable given the fact that the intersystem crossing via the minimum energy crossing point has the lowest barrier in all reaction routes of the triplet state ${ }^{3} \mathrm{CH}_{3} \mathrm{CCHO}$ as shown in the theoretical calculation of this work. The minimum energy crossing point results into fast intersystem crossing, from which nonadiabatic reaction pathways such as $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ are open and play significant roles.
3.3. Comparison with Previous Experiments. Besides the above-discussed CO channel, the calculation reveals the formation mechanisms of a variety of other channels observed by previous crossed molecular beam or LIF experiments ${ }^{17,18}$ as follows.

The H Channels. Our computational results show that H atom can be generated from three pathways: (i) the acetylentic H -atom displacement from IM1; (ii) the methyl H -atom displacement from IM1; (iii) the methyl H-atom displacement from IM2, the methylketene. Channel (i) producing acetylentic H -atom should dominate over channels (ii) and (iii) producing methyl H-atom because the former undergoes a much lower energy barrier, i.e., TS2 compared to TS3 and TS4. This prediction can explain nicely the previous experiments results. Kanofsky et al. ${ }^{17}$ observed that the $\mathrm{H} / \mathrm{D}$ ratios were about 0.1 and 10 for $\mathrm{CH}_{3} \mathrm{CCD}$ and $\mathrm{CD}_{3} \mathrm{CCH}$ as reactants with $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$, respectively. Beroshn et al. ${ }^{18}$ obtained a H/D ratio of $0.770 \pm 0.076$ while determining the sources of H atoms with $\mathrm{CH}_{3} \mathrm{CCD}$ as the reactant. Both experiments agree qualitatively that the acetylentic H -atom
displacement is more advantageous over the methyl H -atom displacement.

The $\boldsymbol{C H}_{3}$ Channels. Kanosfsky et al. determined the source of $\mathrm{CH}_{3}$ product with $\mathrm{CD}_{3} \mathrm{CCH}$ as a reactant, indicating that $\mathrm{CH}_{3}$ is nearly all from the initial methyl group of propyne. ${ }^{17}$ This is consistent with our computational results. The product $\mathrm{CH}_{3}$ can be formed either from the direct displacement of $\mathrm{CH}_{3}$ group from the C2-atom adduct IM5 via TS14 or from the displacement of $\mathrm{CH}_{3}$ from triplet methylketene IM2 via TS8 following the $1,2-\mathrm{H}$ shift of the C 1 -atom adduct IM1 via TS4. Certainly there is no exchange between methyl H and acetylentic H .

The HCO Channel and the $\mathbf{C H}_{2}$ Channel. The calculation shows that HCO is formed on the triplet surface via the displacement of HCO moiety from IM 3 following the rearrangement of IM1 $\left(\mathrm{CH}_{3} \mathrm{CCHO}\right)$ to $\mathrm{IM} 3\left(\mathrm{CH}_{2} \mathrm{CHCO}\right)$ via a 3,2-H shift. While the most feasible pathway forming $\mathrm{CH}_{2}$ is the isomerization of the C 2 -atom addition adducts $\mathrm{IM} 5\left(\mathrm{CH}_{3} \mathrm{COCH}\right)$ to IM6 $\left(\mathrm{CH}_{2} \mathrm{COCH}_{2}\right)$ via a $3,1-\mathrm{H}$ shift with the subsequent release of $\mathrm{CH}_{2}$ and $\mathrm{CH}_{2} \mathrm{CO}$. These two channels occurring via the H -migration followed by decomposition are both energetically and kinetically accessible and thus should account for some major yields of the total reaction. The experimental measurement of Kanofsky et al. ${ }^{17}$ also indicated these two to be among major channels. Furthermore, our calculation indicates that there is no H -atom exchange from methyl to acetylentic end when HCO is displaced from IM3, agreeing well with the observation ${ }^{17}$ that DCO yield had a remarkable advantage over HCO yield for the reaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with $\mathrm{CH}_{3} \mathrm{CCD}$.
To interpret the observation of the $\mathrm{CH}_{2} \mathrm{CO}+\mathrm{CH}_{2}$ channel, Kanofsky suggested a possible mechanism via two steps of H shifts: $\mathrm{O}+\mathrm{CH}_{3} \mathrm{CCH} \rightarrow \mathrm{CH}_{3} \mathrm{C}(\mathrm{O}) \mathrm{CH} \rightarrow \mathrm{CH}_{2}(\mathrm{OH}) \mathrm{CH} \rightarrow$ $\mathrm{CH}_{2} \mathrm{C}(\mathrm{O}) \mathrm{CH}_{2} \rightarrow \mathrm{CH}_{2} \mathrm{CO}+\mathrm{CH}_{2} .{ }^{17}$ In order to validate the feasibility of this route, we also calculated the corresponding pathway, IM5 $\rightarrow$ TS17 $\rightarrow$ IM7 $\rightarrow$ TS18 $\rightarrow$ IM8 $\rightarrow$ TS19 $\rightarrow$ IM6 $\rightarrow$ TS16 $\rightarrow \mathrm{CH}_{2} \mathrm{CO}+\mathrm{CH}_{2}$. As shown in Figure 4 (b), this route turns out to be energetically inaccessible because it involves surmounting TS17 with the energy barrier 13.8 kcal $\mathrm{mol}^{-1}$ above the reactants. Instead, another feasible pathway via $3,1-\mathrm{H}$ shift of IM5 $\left(\mathrm{CH}_{3} \mathrm{COCH}\right)$ to IM6 $\left(\mathrm{CH}_{2} \mathrm{COCH}_{2}\right)$ and the subsequent decomposition of IM6 is revealed in our calculation.

## 4. Conclusions

In brief, the product channels and mechanisms of the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ $+\mathrm{CH}_{3} \mathrm{CCH}$ reaction are investigated both theoretically and experimentally. Theoretically, both of the reaction pathways on the triplet and singlet potential energy surfaces are calculated at the level of $\operatorname{CCSD}(\mathrm{T}) / / \mathrm{B} 3 \mathrm{LYP} / 6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ and the minimum energy crossing point between the two surfaces are located with the Newton-Lagrange method. The theoretical calculations show that the reaction occurs dominantly via the O-addition rather than H -abstraction mechanism. The reaction starts with the O -addition to either of the triple bond carbon atoms forming triplet ketocarbene ${ }^{3} \mathrm{CH}_{3} \mathrm{CCHO}$ or ${ }^{3} \mathrm{CH}_{3} \mathrm{COCH}$ which can undergo decomposition, H -atom migration or intersystem crossing from which a variety of reaction channels become accessible, including the adiabatic channels of $\mathrm{CH}_{3} \mathrm{CCO}+\mathrm{H}\left(\mathrm{CH}_{2} \mathrm{CCHO}\right.$ $+\mathrm{H}), \mathrm{CH}_{3}+\mathrm{HCCO}, \mathrm{CH}_{2} \mathrm{CH}+\mathrm{HCO}, \mathrm{CH}_{2} \mathrm{CO}+\mathrm{CH}_{2}, \mathrm{CH}_{3} \mathrm{CH}$ +CO , and the nonadiabatic channels of $\mathrm{C}_{2} \mathrm{H}_{4}+\mathrm{CO}, \mathrm{C}_{2} \mathrm{H}_{2}+$ $\mathrm{H}_{2}+\mathrm{CO}$, and $\mathrm{H}_{2}+\mathrm{H}_{2} \mathrm{CCCO}$.

Experimentally with TR-FTIR emission spectroscopy, a complete detection of the vibrationally excited product CO up to $v=5$ reveals that the energy released into the CO vibration (i.e., $6.2 \mathrm{kcal} \mathrm{mol}^{-1}$ ) consists with the statistical partition of
the singlet $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{1} \mathrm{~A}\right)+\mathrm{CO}$ channel, but not with the triplet $\mathrm{CH}_{3} \mathrm{CH}+\mathrm{CO}$ or $\mathrm{C}_{2} \mathrm{H}_{4}\left({ }^{3} \mathrm{~A}\right)+\mathrm{CO}$ channel. In combination with the present calculation results, it is concluded that CO arises mainly from the singlet methylketene ( ${ }^{1} \mathrm{CH}_{3} \mathrm{CHCO}$ ) dissociation following the intersystem crossing of the triplet ketocarbene adduct ( $\left.{ }^{3} \mathrm{CH}_{3} \mathrm{CCHO}\right)$. Fast intersystem crossing via the minimum energy crossing point of the triplet and singlet surfaces is shown to play significant roles resulting into nonadiabatic pathways for this reaction. The minimum energy crossing point is located to be $13.9 \mathrm{kcal} \mathrm{mol}^{-1}$ above the triplet ${ }^{3} \mathrm{CH}_{3} \mathrm{CCHO}$, corresponding to the lowest energy barrier among all pathways on the triplet surface.

Overall, the initially formed O-addition ketocarbene adducts $\mathrm{CH}_{3} \mathrm{CCHO}^{*}$ or $\mathrm{CH}_{3} \mathrm{COCH}^{*}$ exhibits full scope of unimolecular decomposition and internal H -atom migration for the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+$ $\mathrm{CH}_{3} \mathrm{CCH}$ reaction. The addition of the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ to either of the carbon atoms of the triple bond is predicted to have almost equal chances because of the nearly identical barrier heights. Following the initial O -addition, three types of H -atom migrations including $1,2-\mathrm{H}$ shift, $3,2-\mathrm{H}$ shift, and $3,1-\mathrm{H}$ shift are found to play equally important roles, from which additional channels become accessible.

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[^1]:    ${ }^{a}$ Calculated at the B3LYP/6-311G(d, p) levels of theory. ${ }^{b}$ Calculated at the CCSD $(\mathrm{T}) / 6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ levels of theory, including ZPE. ${ }^{c}$ The experimental reaction enthalpies at $298 \mathrm{~K}\left(\right.$ in $\mathrm{kcal} \mathrm{mol}^{-1}$ ). The formation enthalpies are taken from refs 32,33 and 34 .

